

# Hydrostatic pressure on $\text{HgBa}_2\text{CaCu}_2\text{O}_{6+\delta}$ and $\text{HgBa}_2\text{Ca}_2\text{Cu}_3\text{O}_{8+\delta}$

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The superconducting transition temperature  $T_c$  and its pressure dependence  $dT_c/dP$  of  $\text{HgBa}_2\text{CaCu}_2\text{O}_{6+\delta}$  (Hg-1212) and  $\text{HgBa}_2\text{Ca}_2\text{Cu}_3\text{O}_{8+\delta}$  (Hg-1223) were measured up to 17 kbar.  $T_c$  increases with pressure approximately linearly for both compounds before oxidation. However, the nonlinearity in the  $T_c$ - $P$  correlation shows up after oxidation in both compounds. For Hg-1212, the average  $dT_c/dP$  decreases as oxygen doping increases, while it increases in Hg-1223. These observations are in conflict with the modified pressure-induced charge-transfer model, but might be attributed to the possible existence of fine electronic structure. These observations suggest that a  $T_c$  much higher than 140 K might be achievable in Hg-1223 by means of higher physical or chemical pressure.

## I. INTRODUCTION

Previous pressure studies on high  $T_c$  cuprates at various oxidation states show that the  $T_c$  changes linearly at pressures between 0 and 20 kbar.<sup>1</sup> This can usually be explained by pressure-induced charge-transfer,<sup>2</sup> i.e.,  $dT_c/dP$  decreases with the carrier concentration  $n$ , which increases with oxidation. Soon after the discovery of the superconducting homologous series  $\text{HgBa}_2\text{Ca}_{n-1}\text{Cu}_n\text{O}_{2n+2+\delta}$  [ $\text{Hg-12}(n-1)n$ ],<sup>3</sup> band-structure calculations showed that the electronic structure of these compounds was strongly affected by the van Hove singularity and evolved with doping rather abnormally,<sup>4</sup> suggesting that their pressure effect  $dT_c/dP$  might be unusual. After samples of these compounds with different oxygen stoichiometry became available, we systematically studied their pressure dependence.

A positive  $dT_c/dP$  was observed in all  $\text{HgBa}_2\text{CaCu}_2\text{O}_{6+\delta}$  (Hg-1212) and  $\text{HgBa}_2\text{Ca}_2\text{Cu}_3\text{O}_{8+\delta}$  (Hg-1223) samples. For the as-prepared Hg-1212 samples with  $T_c \sim 112$  K,  $T_c$  increased linearly with pressure at a rate  $dT_c/dP = 0.22 \pm 0.02$  K/kbar throughout our pressure range. For oxygenated Hg-1212 samples with  $T_c \sim 119$  K,  $T_c$  increases with pressure at a similar rate below  $\sim 4$  kbar, but the rate of increase decreases to  $0.11 \pm 0.02$  K/kbar thereafter. Although the nonlinear  $T_c$ - $P$  correlation is rather unusual, the smaller average  $dT_c/dP$  after oxygenation is consistent with the charge-transfer model,<sup>2</sup> which suggests that the average  $dT_c/dP$  decreases with the carrier concentration  $n$ . The  $dT_c/dP$  of several Hg-1223 samples was also measured. The  $dT_c/dP$  of vacuum-annealed and as-synthesized samples was independent of  $P$ . However, the  $T_c$  vs  $P$  of oxygenated Hg-1223 shows positive curvature above 12 kbar. Unlike that for Hg-1212, the overall  $dT_c/dP$  for Hg-1223 increases with oxidation, which is in direct conflict with the modified charge-transfer model.<sup>2</sup>

## II. EXPERIMENT

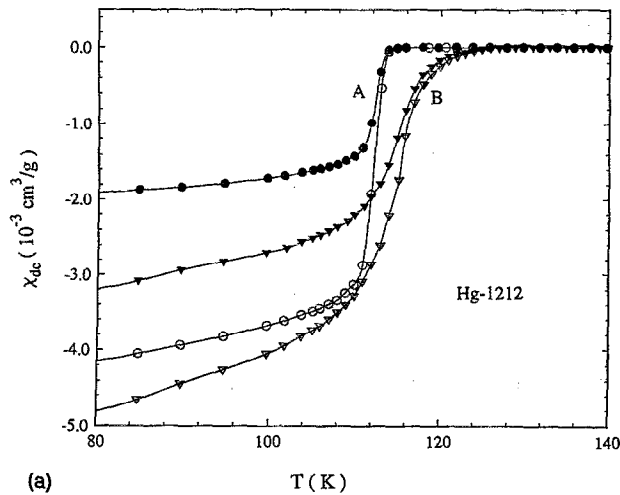
High quality samples were prepared using the controlled vapor/solid reaction (CVSR) technique.<sup>5</sup> The samples are  $\sim 90\%$  pure with  $\sim 10\%$   $\text{CaHgO}_2$  and  $\text{Ba}_2\text{Cu}_3\text{O}_{5+\delta}$  impurities. Structure characterization was carried out by x-ray diffraction using a Rigaku D-MAX/BII powder diffractometer. The dc magnetic susceptibility ( $\chi_{dc}$ ) of samples at ambient

pressure was measured by a quantum design superconducting quantum interference device (SQUID) magnetometer. A standard inductance bridge, operated at 16 Hz in an ac field of  $\sim 5$  Oe peak to peak, was employed to determine the ac magnetic susceptibility ( $\chi_{ac}$ ) under pressures. The standard four-lead measurement was applied to measure resistivity. Hydrostatic pressure up to 18 kbar was generated at room temperature inside a Teflon cup housed in a Be-Cu high pressure clamp,<sup>6</sup> using 3M Fluorinert as the pressure medium. The pressure was determined by a Pb-manometer placed next to the sample. The temperature was measured by an alumel-chromel thermocouple above 30 K and a Ge thermometer below 30 K.

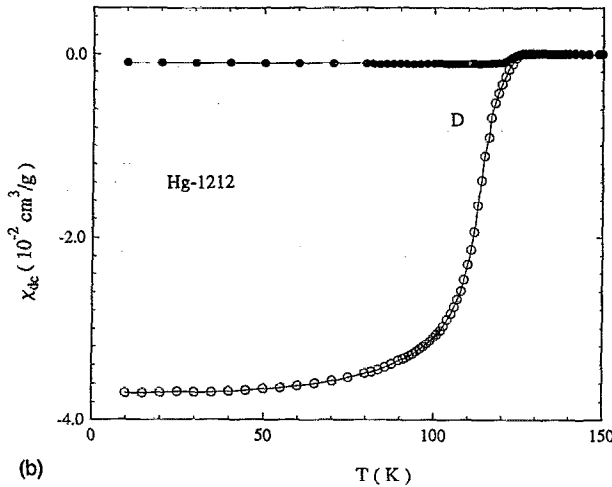
## III. RESULTS AND DISCUSSION

Four samples of Hg-1212 were characterized. Samples A and C were as synthesized with  $T_c \sim 112$  K. Samples B and D were annealed in 1 atm  $\text{O}_2$  at  $300^\circ\text{C}$  for 30 h with  $T_c \sim 119$  K and were determined to be overdoped. The  $\chi_{dc}$  of samples A, B, and D are shown in Fig. 1. The transition is quite sharp and the superconducting volume fraction is large (especially for sample D). The onset and midpoint of  $T_c$  ( $T_{co}$  and  $T_{cm}$ ) as determined by ac-susceptibility measurements at various pressures are shown in Fig. 2. The  $T_{co}$ 's of sample A and C have the same trend as the  $T_{cm}$ 's although they are not shown in the figure. The  $T_{cm}$  of sample D is not well defined due to poor grain coupling and thus not shown.  $T_c$  increases linearly with  $P$  at a rate of  $0.22 \pm 0.02$  K/kbar for samples A and C. For samples B and D,  $T_c$  increases at a rate of  $\sim 0.23$  K/kbar below 4 kbar. The rate fell to  $0.11 \pm 0.03$  K/kbar above 4 kbar. As proposed by the modified charge-transfer model,<sup>2</sup> the overall  $dT_c/dP$  decreases as oxidation increases.

Five samples from two different batches of Hg-1223 were also examined. Sample V was as synthesized with  $T_c \sim 117$  K. Sample W, with  $T_c \sim 79$  K, was obtained by high vacuum annealing at  $450^\circ\text{C}$  for 1 day. Samples X, Y, and Z were annealed in flowing oxygen at  $300^\circ\text{C}$  for 1 day, with  $T_c \sim 135$  K. Figure 3 shows  $T_c$  versus pressure for Hg-1223. The  $T_c$ 's for samples V, X, Y, and Z were measured resistively. Although we only plotted  $T_{cm}$  for these samples, the behavior of  $T_{co}$  and  $T_{cz}$  (zero point of the transition) is similar. A long tail in the superconducting transition was ob-



(a)



(b)

FIG. 1. (a)  $\chi_{dc}$  vs  $T$  for Hg-1212 sample A and B. (b)  $\chi_{dc}$  vs  $T$  for Hg-1212 sample D. Filled symbol: field cooling; open symbol: zero-field cooling.

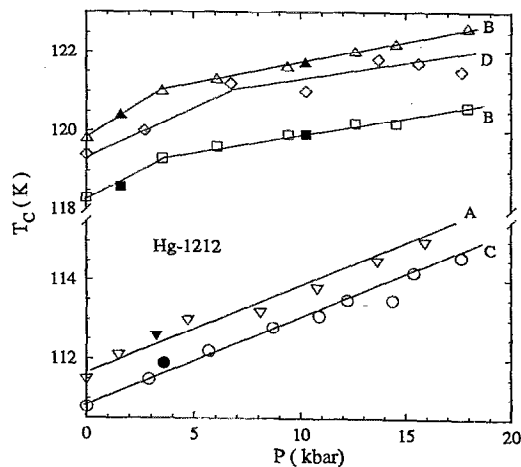


FIG. 2.  $T_c$  vs  $P$  for Hg-1212.  $\Delta$ :  $T_{c0}$  for sample B;  $\diamond$ :  $T_{c0}$  for sample D;  $\square$ :  $T_{cm}$  for sample B;  $\nabla$ :  $T_{cm}$  for sample A;  $\circ$ :  $T_{cm}$  for sample C. Filled symbol obtained on pressure reduction.

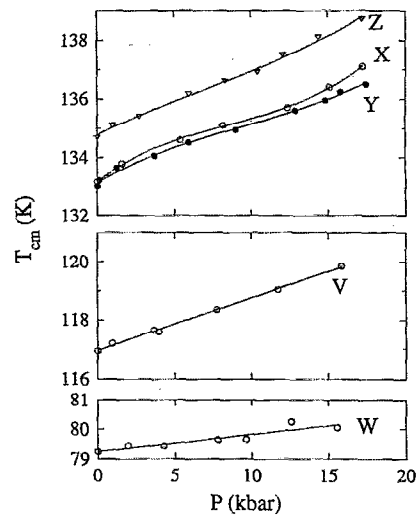


FIG. 3.  $T_{cm}$  vs  $P$  for Hg-1223 samples V, W, X, Y, and Z.

served for the sample W, which might be due to damage caused to the grain boundaries during annealing. Therefore, the  $T_c$  of this sample was measured magnetically and  $T_{cm}$  was obtained from the peak of  $d\chi_{ac}/dT$ . We can see clearly that the average  $dT_c/dP$  increases with oxidation. From thermoelectric measurements, we extracted the carrier concentration  $n$  by an empirical law.<sup>7</sup> Figure 4 shows the average  $dT_c/dP$  vs  $n$  for Hg-1212 and Hg-1223.

Recently, a study<sup>2</sup> of combined effects of pressure and doping on  $(Y,Ca)(Ba,Ca)_2Cu_3O_{7-\delta}$  compounds shows that pressure-induced charge-transfer alone cannot explain the observed results and an "intrinsic" pressure effect on  $T_c$  is included. This term,  $dT_{c,max}/dP$ , is assumed to be independent of  $P$  and  $n$ . Thus, the empirical modified pressure-induced charge-transfer model can be written as

$$T_c(n, P) = T_{c,max}(n_{optimal}, 0) + (dT_{c,max}/dP) \times P + A \times [(n - n_{optimal}) + (dn/dP) \times P]^2,$$

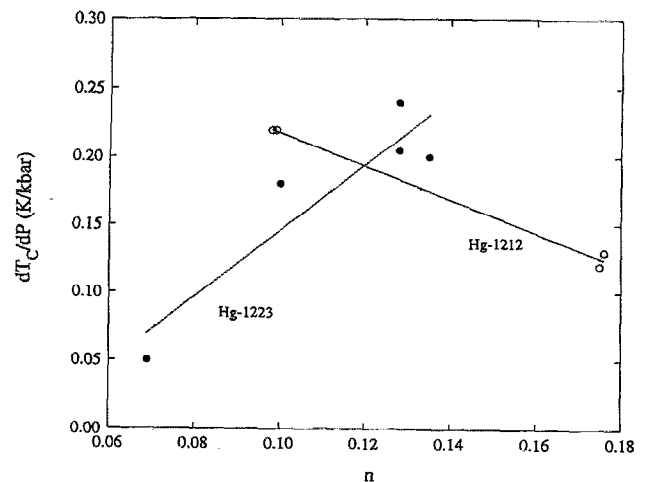


FIG. 4.  $dT_c/dP$  vs  $n$  for Hg-1212 and Hg-1223.

where  $A$  is a negative compound-dependent constant. By ignoring higher order contributions, the pressure effect on  $T_c$  becomes

$$dT_c/dP = dT_{c,\max}/dP + 4A \times dn/dP \times (n - n_{\text{optimal}}),$$

which should decrease with  $n$  (oxidation), since  $dn/dP > 0$  (which has been demonstrated by the negative  $d\rho/dP$ , positive  $dR_H/dP$ ;<sup>8</sup> where  $\rho$  is normal state resistivity and  $R_H$  is the Hall coefficient). However, the average  $dT_c/dP$  increases with  $n$  in Hg-1223, despite the negative  $d\rho/dP$  over the whole  $P$  range. This is in direct conflict with the model<sup>2</sup> and suggests that  $dT_{c,\max}/dP$  should be  $P$  and/or  $n$  dependent. The ever increasing  $dT_c/dP$  with  $P$  after oxidation (cf. Fig. 3) suggests that the  $T_c$  of Hg-1223 can be enhanced further by applying pressures up to 500 kbar.<sup>9</sup>

In summary, we investigated the pressure effect on samples of Hg-1212 and Hg-1223 with different oxygen contents. We found that the pressure effect becomes rather nonlinear after oxidation. Although a modified charge-transfer model can explain the pressure effect on Hg-1212, it is in direct conflict with the observation on Hg-1223. The positive  $dT_c/dP$  for all the samples suggests that high  $T_c$  is optimistic.

## ACKNOWLEDGMENTS

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- <sup>1</sup>J. S. Schilling and S. Klotz, in *Physical Properties of High Temperature Superconductors III*, edited by D. M. Ginsberg (World Scientific, Singapore, 1992), p. 84.
- <sup>2</sup>J. J. Neumeier and H. A. Zimmermann, *Phys. Rev. B* **47**, 8386 (1993).
- <sup>3</sup>S. N. Putilin, E. V. Antipov, O. Chmaissem, and M. Marezio, *Nature* **362**, 226 (1993).
- <sup>4</sup>D. L. Novikov, V. A. Gubanov, and A. J. Freeman, *Physica C* **210**, 301 (1993); D. L. Novikov and A. J. Freeman, *Physica C* **216**, 273 (1993).
- <sup>5</sup>R. L. Meng, L. Beauvais, X. N. Zhang, Y. Y. Sun, J. Kulik, Y. Y. Xue, and C. W. Chu, *Physica C* **216**, 21 (1993).
- <sup>6</sup>C. W. Chu, *Phys. Rev. Lett.* **33**, 1283 (1974).
- <sup>7</sup>S. D. Obertelli, J. R. Cooper, and J. L. Tallon, *Phys. Rev. B* **46**, 14928 (1992).
- <sup>8</sup>C. Murayama, Y. Iye, T. Enomoto, A. Fukushima, N. Môri, Y. Yamada, and T. Matsumoto, *Physica C* **185-189**, 1293 (1991).
- <sup>9</sup>L. Gao, Y. Y. Xue, F. Chen, Z. J. Huang, R. L. Meng, and C. W. Chu, *Nature* **365**, 323 (1993).